

# Summary of Beating and Resonance

## The Equations

We consider oscillators with no damping and periodic forcing, and we think of the forcing as having an adjustable period (we can change  $\omega$ , but  $b$  is fixed).

$$y'' + b^2y = \cos(\omega t) \quad y(0) = 0, \quad y'(0) = 0$$

Initially, we assume  $\omega \neq b$ . The full solution to the IVP was then:

$$y(t) = \frac{1}{b^2 - \omega^2}(\cos(\omega t) - \cos(bt))$$

It is possible to write this expression as:

$$y(t) = \frac{1}{b^2 - \omega^2} \left( -2 \sin\left(\frac{\omega + b}{2}t\right) \sin\left(\frac{\omega - b}{2}t\right) \right)$$

And when  $\omega = b$ , the solution becomes:  $y(t) = \frac{1}{2b}t \sin(bt)$

## The Analysis

When considering solutions to the undamped mass-spring model with periodic forcing, we found that, as we vary the frequency of the forcing function, as it gets close to the natural (homogeneous) solution, then **beating** will begin to occur. As the forcing frequency matches the natural frequency, the forced response will “blow up”, which is **resonance**.

We also found that the beat period (for the “slow” function) is half the period of the slower sine:

$$\frac{1}{2} \cdot \frac{2\pi}{\frac{\omega-b}{2}} = \frac{\pi}{\frac{\omega-b}{2}} = \frac{2\pi}{\omega - b}$$

and the beat amplitude was:

$$A = \frac{2}{|b^2 - \omega^2|}$$

The period of the “fast” beat is the period of  $\sin((\omega + b)t/2)$ .

These things together say that, as  $\omega \rightarrow b$ , the amplitude and the period of the beats get larger, and larger and larger- Finally blowing up when they are equal.

For example, if  $\omega = 2.17$  and  $b = 2$ , then we'll get beating, with:

$$\text{Period of one beat} \approx \frac{2\pi}{0.17} \approx 11.76\pi$$

and the amplitude of the beat is:  $\frac{2}{|b^2 - \omega^2|} \approx \frac{2}{0.71} \approx 2.82$ .

Something to consider: In nature, there is never the situation where there is no damping at all, but there may be damping that is so small as to be basically 0.