Extra Practice: Exercise 4.4.3

Consider a slightly damped vibrating string that satisfies:

$$\rho_0 u_{tt} = T_0 u_{xx} - \beta u_t$$

1. Briefly explain why $\beta > 0$.

Extra HINT: Look at the how the acceleration and u are related by considering

 $\rho_0 u_{tt} = -\beta u_t$

SOLUTION: We see here that with $\beta > 0$, the force of acceleration is acting in the opposite direction as the velocity, so this is *friction*.

2. Assume that ρ_0 , T_0 and β are constants, and determine the solution by separation of variables that satisfy the given conditions:

BCs
$$u(0,t) = 0$$
 $u(L,t) = 0$
ICs $u(x,0) = f(x)$ $u_t(x,0) = g(x)$

And assume that β is small, $\beta^2 < 4\pi^2 \rho_0 T_0/L^2$.

SOLUTION: Let u = XT as usual, and substitute into the PDE to get

$$\rho_0 X T'' = T_0 X'' T - \beta X T' \quad \Rightarrow \quad \rho_0 X T'' + \beta X T' = T_0 X'' T$$

With the boundary conditions, it might be easiest to keep the constants with T. To do that, divide both sides by T_0XT to get:

$$\frac{\rho_0 T'' + \beta T'}{T_0 T} = \frac{X''}{X} = -\lambda$$

Analyzing the spatial ODE first, we get our familiar BVP

$$X'' + \lambda X = 0,$$
 $X(0) = 0$ $X(L) = 0$

so that $\lambda = 0$, $\lambda < 0$ lead us to the trivial solution, and the eigenvalues and eigenfunctions are:

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \qquad X_n(x) = \sin\left(\frac{n\pi}{L}x\right)$$

Now we solve the time-dependent ODE

$$\rho_0 T'' + \beta T' = -\lambda T_0 T \quad \Rightarrow \quad \rho_0 T'' + \beta T' + \lambda T_0 T = 0$$

The characteristic equation is, and solve using the quadratic formula:

$$\rho_0 r^2 + \beta r + \lambda T_0 = 0 \quad \Rightarrow \quad r = \frac{-\beta \pm \sqrt{\beta^2 - 4\lambda T_0 \rho_0}}{2\rho_0}$$

Using the assumption (and the definition of λ_n), we can show that the roots are complex:

$$\beta^2 < \frac{4\pi^2 \rho_0 T_0}{L^2} < 4 \frac{n^2 \pi^2}{L^2} \rho_0 T_0 \quad \text{for } n = 1, 2, 3, \cdots$$

Therefore, the discriminant is negative, and the roots are complex. We also don't want to get hung up in the notation, so let's make a couple of substitutions:

$$r = -\frac{\beta}{2\rho_0} \pm \frac{\sqrt{4\lambda_n T_0 \rho_0 - \beta^2}}{2\rho_0} i = \gamma \pm \omega_n \eta_0$$

(note that ω depends on λ_n , thus the ω_n notation). Using this substitution, the solutions in time are the following, and for future reference, the derivative is given as well:

$$T_n(t) = e^{\gamma t} \left(A_n \cos(\omega_n t) + B_n \sin(\omega_n t) \right)$$

with

$$\Gamma_n'(t) = \gamma e^{\gamma t} \left(A_n \cos(\omega_n t) + B_n \sin(\omega_n t) \right) + e^{\gamma t} \left(-A_n \omega_n \sin(\omega_n t) + B_n \omega_n \cos(\omega_n t) \right)$$

Taking the superposition:

r

$$u(x,t) = \sum_{n=1}^{\infty} e^{\gamma t} (A_n \cos(\omega_n t) + B_n \sin(\omega_n t)) \sin\left(\frac{n\pi}{L} x\right)$$

To find A_n, B_n , we use the initial position and velocity:

$$u(x,0) = f(x) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi}{L}x\right) \quad \Rightarrow \quad A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi}{L}x\right) \, dx$$

And differentiating, we get the following (the derivative was computed earlier):

$$u_t(x,0) = g(x) = \sum_{n=1}^{\infty} T'_n(0) \sin\left(\frac{n\pi}{L}x\right) = \sum_{n=1}^{\infty} \left(\gamma A_n + \omega_n B_n\right) \sin\left(\frac{n\pi}{L}x\right)$$

Notice that the A_n have already been computed. With this, we can solve for the B_n :

$$(\gamma A_n + \omega_n B_n) = \frac{2}{L} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx$$

with

$$B_n = -\frac{\gamma}{\omega_n} A_n + \frac{2}{\omega_n L} \int_0^L g(x) \sin\left(\frac{n\pi}{L}x\right) dx$$

A little commentary: As nasty as this may look, adding friction (β) ended up just multiplying our previous solution(s) by

$$e^{\gamma t} = e^{\frac{-\beta}{2\rho_0}t}$$

which dampens out the oscillations. Adding the dampening also affected the natural frequencies as well, but this was expected- the same thing happened in Math 244 in our analysis of the mass-spring system:

$$mu'' + \gamma u' + ku = 0$$

where γ was friction, k was the spring constant, m was mass.